# The Crystal Structure of $\mathbf{N a}_{2} \mathrm{Ti}_{7} \mathbf{O}_{15}$, an Ordered Intergrowth of $\mathrm{Na}_{2} \mathbf{T i}_{6} \mathbf{O}_{13}$ and ' $\mathbf{N a}_{2} \mathbf{T i}_{8} \mathbf{O}_{17}$ ' 

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#### Abstract

Crystals of $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$ were isolated from a specimen of composition $\mathrm{Na}_{2} \mathrm{O} \cdot \mathrm{Al}_{2} \mathrm{O}_{3} .2 \mathrm{TiO}_{2}$ heated in a sealed platinum capsule at $1300^{\circ} \mathrm{C}$. Their symmetry is monoclinic, space group $C 2 / \mathrm{m}$, with the unitcell dimensions $a=14.9, b=3.74, c=20.9 \AA, \beta=96.5^{\circ}$. The structure, deduced from the Patterson projection $P(u, w)$ by image-seeking and refined by three-dimensional difference Fourier methods, contains two kinds of condensed Ti-O octahedral group. One is three octahedra wide, joined by multiple edge-sharing to other units of three and forming an endless zigzag ribbon, whilst the other is four wide, but similarly joined. The groups unite with each other by corner-sharing to form a series of open tunnels containing the Na ions. The structure can be regarded as an ordered intergrowth, at the unit-cell level, of sodium hexatitanate $\mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}$, and a postulated sodium octatitanate $\mathrm{Na}_{2} \mathrm{Ti}_{8} \mathrm{O}_{17}$.


## Introduction

Recently we described the phase $\mathrm{NaTi}_{2} \mathrm{Al}_{5} \mathrm{O}_{12}$ which resulted from an effort to synthesize a compound with the composition $\mathrm{NaAlTi}_{2} \mathrm{O}_{6}$ (Mumme \& Wadsley, 1967). The two metals Ti and Al in the structure proved to be ordered, while the coordination of Al was of two kinds, octahedral and tetrahedral. The association of these same elements in a sodium-containing mixed oxide was reported for non-stoichiometric $\mathrm{Na}_{x}\left(\mathrm{Al}_{x} \mathrm{Ti}_{1}-x\right) \mathrm{O}_{2}, x$ approximately $0 \cdot 2$ (Bayer \& Hoffman, 1965), but in this case the role of $\mathrm{Al}^{3+}$ was to replace $\mathrm{Ti}^{4+}$ by an amount necessary to maintain the balance of charge. The compound has the $\mathrm{Na}_{x} \mathrm{TiO}_{2}$ structure (Andersson \& Wadsley, 1962b).

Attempts in this laboratory to prepare $\mathrm{NaAlTiO}_{4}$, a possible member of an isomorphous group $\mathrm{Na} A^{3+} B^{4+} \mathrm{O}_{4}$ with the $\mathrm{CaFe}_{2} \mathrm{O}_{4}$ structure (Reid, Wadsley \& Sienko, 1967) yielded a heterogeneous matrix containing a small number of colourless needle-like crystals. These were shown by the present structure analysis to be essentially $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$, and not an aluminotitanate as originally supposed. This compound proved to be one of a series of structurally related alkali titanates, the already known members being $\mathrm{Na}_{2} \mathrm{Ti}_{3} \mathrm{O}_{7}$ (Andersson \& Wadsley, 1961) and $\mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}$ (Andersson \& Wadsley, 1962a).

## Experimental

Sodium oxalate, aluminum oxide and titanium dioxide (anatase) in the molar ratios 1:1:2 were intimately mixed, prefired at $900^{\circ} \mathrm{C}$ to remove carbon dioxide, reground, and heated for 24 hours at $1300^{\circ} \mathrm{C}$ in a sealed platinum capsule. The small needles were the only crystalline material large enough to manipulate, and no attempt was made to identify the remaining material. The Guinier powder pattern contained few lines, many having indexing ambiguities. The lattice parameters in Table 1 were therefore evaluated from single-crystal films; no greater accuracy than $\pm 0 \cdot 1 \AA$ for $a$ and $c$ could be obtained.

Table 1. Crystallographic constants for $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$
Symmetry: monoclinic
Unit-cell dimensions: $a=14 \cdot 9 \pm 0 \cdot 1 \AA$
$b=3.74 \pm 0.01$
$c=20.9 \pm 0.1$
$\beta=96.5 \pm 0.5$
Systematically
absent reflexions: $h k l$ with $h+k \neq 2 n$
Possible space groups: $C 2 / m$ (no. 12), $C m$ (no. 8), $C 2$ (no. 5)

$$
\begin{aligned}
& D_{x}=3.57 \mathrm{~g} . \mathrm{cm}^{-3} \\
& Z=4
\end{aligned}
$$

$\mu=322 \mathrm{~cm}^{-1}$
Intensities were collected from a crystal measuring $0.05 \times 0.02 \times 0.01 \mathrm{~mm}$ oriented for rotation about the $b$ axis, the direction of elongation. The levels $h k l$, $k=0,1$ and 2 were recorded on multiple film packs by the integrating Weissenberg method with copper $K \alpha$ radiation. Intensities were estimated visually with a standard scale, and specimen absorption was neglected. Only 167 reflexions, a small fraction of those in the copper sphere, were strong enough to be measurable. Scattering curves for the neutral atoms Ti and Na were taken from International Tables for $X$-ray Crystallography (1962) while Suzuki's (1960) values for $\mathrm{O}^{2-}$ were used for oxygen. Structure amplitudes were approximately scaled by Wilson's method, and the individual scale factors adjusted by comparing $F_{o}$ and $F_{c}$ before each successive refinement cycle. All calculations were made on the Division's Elliott 803 computer with the programming system of Daly, Stephens \& Wheatley (1963), kindly made available by P.J. Wheatley and J.J.Daly.

## Structure determination

The Patterson projection $P(u, w)$, consisted of parallel rows of well-defined peaks extending at angles of $80^{\circ}$ from the positive direction of the $a$ axis, and subdivided into a nearly rectilinear sub-lattice repeating itself fourteen times along $c$ (Fig. 1). It could safely be assumed that the $b$-axis length, $3.74 \AA$, corresponded
to an octahedral body diagonal. Tentative models of the structure in cases such as this can be rapidly assessed by drawing an orthogonal grid of interconnected and overlapping squares each representing an ideal octahedron in projection, the centre being the metal together with the under- and overlying oxygen atoms, and the corners the remaining oxygen atoms. The orientation of the grid is determined by the relationship of the sub-cell to the unit cell, and the problem is one of finding a particular arrangement of edge- and corner-shared octahedra with the correct unit-cell size and space group, with suitable positions for the sodium atoms and with a balanced chemical composition.

By transferring the origin of a duplicate transparent map of the Patterson function to the strong peak marked $A$ in Fig. 1, the major vector coincidences were represented on the grid as octahedra and not just the individual atoms. Separations of $3 \cdot 0-3 \cdot 1 \AA$ were assumed to arise from octahedra sharing edges at the same level perpendicular to the $b$ axis, while $2 \cdot 5-2 \cdot 6 \AA$ corresponded to edges shared at different levels. In the latter case the squares on the model overlapped. This simple image-seeking procedure, which we have described in detail elsewhere (Gatehouse \& Wadsley, 1964; Andersson, Mumme \& Wadsley, 1966) led to a partial solution of the problem. The structure seemed to contain two identical but independent groups of three octahedra sharing edges in a plane and repeating $b / 2$ downwards by additional edge-sharing, thereby forming two zigzag ribbons extending along the axis of symmetry. At this stage the manner of joining the groups could not be ascertained, nor could realistic positions for sodium be recognized. An electron-density projection on to ( 010 ), however, revealed an additional metal atom. Its position increased the size of one zigzag assemblage to four repetitive octahedra, so that the asymmetric unit contained seven metals - excluding sodium - a conclusion fully confirmed by the 14 -fold periodicity in the Patterson projection.

This structural model was completed by placing two sodium atoms in fourfold positions in the tunnels formed by the junctions of the groups, giving the formula $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$. Three-dimensional electron density sections confirmed this general arrangement, but sev-


Fig.1. Projection of the Patterson function, $P(u, w)$. The strong vector $A$ was the origin selected for image-seeking by superposition.
eral additional subsidiary maxima appeared for which there was no meaningful explanation. These were subsequently eliminated by refinement procedures, and evidently resulted from series termination. A number of iterative three-dimensional difference syntheses were computed, and least-squares refinement was used at the same time to evaluate standard deviations as well as individual isotropic temperature factors. A persistent small peak at the origin of the difference syntheses was identified as a sodium atom in the point position $2(a)$ of $C 2 / m$, and because of this the remaining six sodium atoms in the unit cell were distributed over the two fourfold sites originally chosen. Refinement would not go beyond an $R$ index of $12 \cdot 2 \%$. Final atomic parameters are listed in Table 2, where the standard deviations and temperature factors are averaged out for atoms of the same kind. A comparison of observed and calculated structure factors is given in Table 3. The calculated $F$ values for all reflexions observed zero were all very small and are omitted. The structure is illustrated in Fig. 2 as a ball-and-spoke drawing and as an idealized octahedral model by Fig. 3(b).

Table 2. Fractional atomic parameters for $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$
Space group $C 2 / m$. E.s.d.'s in parentheses.

|  | Point position | $x$ | $y$ | $z$ | $B$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Na}(1)$ | 2(a) | 0 | 0 |  | $2 \cdot 4 \AA^{2}$ |
| ${ }_{4}^{3} \mathrm{Na}(2)$ | 4(i) | 0.962 (6) | 0 | 0.395 (4) | $2 \cdot 4$ |
| ${ }_{4}^{3} \mathrm{Na}(3)$ | 4(i) | 0.069 (6) | 0 | $0 \cdot 183$ (4) | $2 \cdot 4$ |
| $\mathrm{Ti}(1)$ | 4(i) | 0.6965 (27) | 0 | 0.0430 (18) | $1 \cdot 1$ |
| Ti(2) | 4(i) | 0.7477 (27) | 0 | $0 \cdot 1863$ (18) | $1 \cdot 1$ |
| $\mathrm{Ti}(3)$ | 4(i) | 0.6237 (27) | 0 | 0.3279 (18) | $1 \cdot 1$ |
| Ti(4) | 4(i) | 0.6718 (27) | 0 | 0.4717 (18) | $1 \cdot 1$ |
| Ti(5) | 4(i) | $0 \cdot 3520$ (27) | 0 | 0.0996 (18) | $1 \cdot 1$ |
| Ti(6) | 4(i) | 0.4007 (27) | 0 | $0 \cdot 2439$ (18) | $1 \cdot 1$ |
| Ti(7) | 4(i) | $0 \cdot 2775$ (27) | 0 | 0.3852 (18) | $1 \cdot 1$ |
| O(1) | 4(i) | 0.646 (7) | 0 | $0 \cdot 119$ (5) | 1.5 |
| O(2) | 4(i) | 0.686 (7) | 0 | $0 \cdot 250$ (5) | 1.5 |
| O(3) | $4(i)$ | 0.517 (7) | 0 | $0 \cdot 280$ (5) | 1.5 |
| O(4) | $4(i)$ | 0.576 (7) | 0 | $0 \cdot 412$ (5) | 1.5 |
| O(5) | 4(i) | $0 \cdot 185$ (7) | 0 | $0 \cdot 472$ (5) | 1.5 |
| O(6) | 4(i) | $0 \cdot 137$ (7) | 0 | $0 \cdot 328$ (5) | 1.5 |
| O(7) | 4(i) | 0.380 (7) | 0 | 0.449 (5) | 1.5 |
| O(8) | 4(i) | 0.332 (7) | 0 | 0.313 (5) | 1.5 |
| O(9) | $4(i)$ | 0.454 (7) | 0 | $0 \cdot 155$ (5) | 1.5 |
| O(10) | $4(i)$ | 0.402 (7) | 0 | 0.018 (5) | 1.5 |
| O(11) | 4(i) | 0.831 (7) | 0 | $0 \cdot 100$ (5) | $1 \cdot 5$ |
| O(12) | 4(i) | 0.892 (7) | 0 | $0 \cdot 244$ (5) | 1.5 |
| $\mathrm{O}(13)$ | 4(i) | 0.752 (7) | 0 | 0.386 (5) | 1.5 |
| O (14) | $4(i)$ | $0 \cdot 216$ (7) | 0 | $0 \cdot 042$ (5) | 1.5 |
| O(15) | 4(i) | $0 \cdot 266$ (7) | 0 | $0 \cdot 186$ (5) | 1.5 |

Despite high standard deviations, the interatomic distances for the titanium-oxygen octahedra listed in Table 4 show the spread of values normal for low symmetry titanate structures. The environment of sodium is of two kinds. $\mathrm{Na}(2)$ has two bonds each of 2.56 , $2.52,3.22$ and $3.07 \AA$ while $\mathrm{Na}(3)$ has two each of $2 \cdot 64,2 \cdot 56,2 \cdot 93$ and $2 \cdot 82 \AA$ (all $\pm 0 \cdot 13 \AA$ ). Both atoms are displaced from the centres of distorted cubes of oxygen atoms (Fig.2) in exactly the same way as in the hexatitanates $\mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}$ (Andersson \& Wadsley, $1962 a$ ) and $\mathrm{K}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}(\mathrm{Cid}$-Dresdner \& Buerger, 1962).
$\mathrm{Na}(1)$ occupying the twofold position at the origin, on the other hand, is at the centre of a square of oxygens [the atoms $\mathrm{O}(4)$ ], with four equal bonds $2 \cdot 43 \pm 0 \cdot 2 \AA$, which is much the same as Na in the bronze $\mathrm{Na}_{x} \mathrm{Ti}_{4} \mathrm{O}_{8}$ (Andersson \& Wadsley, 1962b).

## Discussion

Problems of the present kind pose a number of difficulties centring around the small percentage of reflexions of measurable intensity which are recorded, despite the large number of atoms in the asymmetric unit. This scarcity of data is caused by the quasi-orthogonal sub-structure dominating the reciprocal lattice. The 'superlattice' reflexions appear sporadically instead of falling close to a sub-structure spot, and the correct assignment of indices, particularly at high angles, is a matter of considerable initial uncertainty. Light atoms in special positions such as $\mathrm{Na}(1)$, contributing equally to all reflexions, may go unnoticed in an electron density map containing only a small fraction of the data. Even a heavy atom under similar circumstances appears deceptively small (Roth \& Wadsley, 1965). Difference syntheses are perhaps the best way of attacking these problems, particularly if appraised in conjunction with temperature factors, which are sensitive to incorrectly weighted scattering curves.

Even so, the distribution of sodium in $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$ must be considered indecisive, as any number of combinations of occupancy and thermal motion are possible. The alkali metals in tunnel compounds behave erratically, and for two members of the class, $\mathrm{NaTi}_{2} \mathrm{Al}_{5} \mathrm{O}_{12}$ (Mumme \& Wadsley, 1967) and $\mathrm{Na}_{0.9} \mathrm{Fe}_{0.9} \mathrm{Ti}_{1.1} \mathrm{O}_{4}$ (Mumme \& Reid, 1968), where in both cases the diffraction data were more complete, the temperature factors in the tunnel directions were very large. This can be attributed to the freedom of these atoms to vibrate relatively unhindered in this particular direction. Nevertheless, in the present case it is conceivable that all of the sodium positions might be completely occupied if aluminum substitutes isomorphously for titanium, when the formula is $\mathrm{Na}_{2 \cdot 5}\left(\mathrm{Ti}_{6 \cdot 5} \mathrm{Al}_{0 \cdot 5}\right) \mathrm{O}_{17}$.

The structural implications of this compound are more significant than its exact stoichiometry. Fig.3(a)


Fig. 2. Structure of $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$. Smallest circles Ti, largest O , intermediate (stippled) Na .

Table 3. Observed and calculated structure factors

Table 4. Interatomic distances for Ti octahedra in $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$

|  |  | Distances ( $\AA$ )order as in column 2) |  | O-O distances ( $\AA$ ) |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Metal | Bonded oxygen atoms* |  | Average | Max. | Min. | Average |
| Ti(1) | 14 (2); 1, 11, 10', 14' | 1.89 (2); 1.83, 2.21, 1.83, 2.32 | 1.99 | 2.97 | 2.73 | $2 \cdot 80$ |
| Ti(2) | 15 (2); 1, 2, 12, 11 | 1.89 (2); 1.95, $1.70,2.34,2.30$ | 2.01 | 3.09 | 2.66 | $2 \cdot 83$ |
| $\mathrm{Ti}(3)$ | 6 (2); 2, 3, 4, 13 | 1.88 (2); 1.96, $1.78,1.97,2.15$ | 1.93 | 2.90 | 2.64 | 2.73 |
| Ti(4) | 5 (2); 4, 7', 5', 13 | 1.88 (2); 1.79, 1.91, 2.28, 2.26 | 2.00 | $3 \cdot 00$ | 2.69 | $2 \cdot 81$ |
| Ti(5) | 11 (2); 9, 10, 14, 15 | 1.90 (2); 1.80, 1.94, 2.24, 2.33 | 2.01 | $3 \cdot 02$ | 2.73 | $2 \cdot 83$ |
| Ti(6) | 12 (2); 8, 3', 9, 15 | 1.87 (2); 1.86, 1.81, 2.10, 2.22 | 1.95 | 2.95 | 2.58 | 2.76 |
| Ti(7) | 13 (2); 5, 7, 8, 6 | 1.91 (2); 2.40, 1.91, 1.79, 2.29 | 2.03 | 3.01 | 2.73 | $2 \cdot 85$ |

[^0]
(a)

(b)

(c)

Fig. 3. Idealized structures of monoclinic sodium titanates in projection on to (010). The squares are Ti-O octahedra; the heavy and light outlines represent the two levels in the unit cells, which are outlined. Na atoms are stippled circles. (a) $\mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}$ (Andersson \& Wadsley, 1962a). The octahedra join, by sharing edges, into groups of three. (b) $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$, with alternating groups of three and four octahedra sharing edges. The sodium positions not at the origin and side centre are incompletely filled. (c) Hypothetical $\mathrm{Na}_{2} \mathrm{Ti}_{8} \mathrm{O}_{17}$, with groups of four octahedra. The distribution of Na is unknown. (b) is a hybrid or intergrowth of (a) and (c).
and (b) show its relationship to sodium hexatitanate $\mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}$ (Andersson \& Wadsley, 1962a), which is built up from zigzag units of octahedra, three wide, and joined in an identical manner to form a threedimensional framework structure. We can also visualze the octatitanate $\mathrm{Na}_{2} \mathrm{Ti}_{8} \mathrm{O}_{17}$ [or perhaps $\mathrm{Na}_{3}\left(\mathrm{AlTi}_{7}\right) \mathrm{O}_{17}$ ] in Fig.3(c), identical with $\mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13}$ except for the size of the octahedral groups which are four wide instead of three. $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$ containing both groups, three and four wide alternating regularly, is thus a hybrid, or an ordered intergrowth at the unit-cell level of these two structures. A number of examples of this particular phenomenon have been reported; for example, the pseudobinary mixed oxide $\mathrm{W}_{4} \mathrm{Nb}_{26} \mathrm{O}_{77}$ contains alternating segments of the two phases $\mathrm{WNb}_{12} \mathrm{O}_{33}$ and $\mathrm{W}_{3} \mathrm{Nb}_{14} \mathrm{O}_{44}$ (Andersson, Mumme \& Wadsley, 1966).

Up to the present, however, the octatitanate has not been prepared. We have no doubt that a compound with its predicted structure can be made, but these titanates $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$ and $\mathrm{Na}_{2} \mathrm{Ti}_{8} \mathrm{O}_{17}$ do not form by the simple expedient of heating the weighed and mixed oxides in a sealed system. The readily made sodium hexatitanate could be expected to decompose in the solid state according to the equation

$$
7 \mathrm{Na}_{2} \mathrm{Ti}_{6} \mathrm{O}_{13} \rightarrow 6 \mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}+\mathrm{Na}_{2} \mathrm{O}
$$

provided it can be heated to a sufficiently high temperature without fusing. The decomposition of one phase to another one of like structure and composition occurs by the cooperative diffusion of non-metals as well as metal ions within the solid (Andersson \& Wadsley, 1966). The formation of $\mathrm{Na}_{2} \mathrm{Ti}_{7} \mathrm{O}_{15}$ from the composition $\mathrm{Na}_{2} \mathrm{O} . \mathrm{Al}_{2} \mathrm{O}_{3} .2 \mathrm{TiO}_{2}$ (i.e. ' $\mathrm{NaAlTiO}{ }_{4}$ ') may simply be due to $\mathrm{Al}_{2} \mathrm{O}_{3}$ altering the phase diagram so as to permit such a diffusion reaction to take place.

The structures of the tri-, hexa- and heptatitanates afford additional examples of two working rules found to be useful in crystal chemical studies of lattice compounds of the transition metals, notably the oxides of $\mathrm{Ti}, \mathrm{V}, \mathrm{Nb}, \mathrm{Mo}$, Ta and W (Wadsley, 1967). The threeand four-member octahedral groups illustrate the first, that an octahedron, centred around a transition metal ion in its maximum valency state, is attracted towards similar octahedra and unites with them so as to have the maximum edge (or face) sharing compatible with the stoichiometry of the phase in which it occurs. The second states that these groups replace the single octahedron as the basic unit of structure, and can join by peripheral corner- or edge-sharing with the same or with different groups to form three-dimensional networks, which may include another kind of metal ion at suitable interstices.

Finally, these titanates may be regarded as derivatives of a parent structure characterized by the ideal stoichiometry $\mathrm{M}^{+} A^{3+} \mathrm{TiO}_{4}, \mathrm{M}^{+}$being an alkali metal $\mathrm{K}, \mathrm{Rb}$ or Cs . The relationships with the alkali titanates will be discussed in a forthcoming publication (Mumme, Reid \& Wadsley, 1968).

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# The Crystal Structure of the 1:2 Copper(II) Nitrate-Methyl Cyanide Complex 

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The structure of $\mathrm{Cu}\left(\mathrm{NO}_{3}\right)_{2} .2 \mathrm{CH}_{3} \mathrm{CN}$ has been solved by the heavy-atom method and refined to $R=0.15$ by least-squares analysis based on 292 partial three-dimensional data. The complex crystallizes in the monoclinic space group $P 2_{1} / c$ with unit-cell dimensions $a=8.81, b=7.36, c=14.73 \AA, \beta=94^{\circ}$, and $Z=4$. Each copper atom in the structure is in a (4+1) tetragonal pyramidal environment and half the nitrate groups act as bridging ligands, linking tetragonal pyramids related by the twofold screw axes, to form infinite chains parallel to the $b$ crystallographic axis. The remaining nitrate groups are unidentate, forming only one bond to the copper atom. Methyl cyanide molecules, coordinated to copper through their nitrogen atoms, occupy trans positions in the copper square plane. The structure is compared with the structures of $\alpha$-cupric nitrate and its complex with nitromethane.

## Introduction

The crystal structure of the $1: 2$ complex between anhydrous cupric nitrate and methyl cyanide has been determined in order to compare the structure of $\alpha$ cupric nitrate (Wallwork \& Addison, 1965) with the crystal structures of its complexes with organic ligands. It would be of interest to determine any changes in metal-nitrate bonding which may occur on complex formation and to investigate the effect of the organic ligands on the copper coordination sphere.

## Experimental

Anhydrous cupric nitrate, prepared by the method of Addison \& Hathaway (1958), was dissolved in freshly distilled methyl cyanide at $40^{\circ} \mathrm{C}$ until a saturated solution was obtained. Deep blue crystals of the methyl cyanide adduct were deposited on standing and analysis

[^1]for copper content confirmed the $1: 2$ composition. Because of the deliquescent nature of the complex, suitable crystals were sealed in thin-walled Pyrex capillary tubes in a dry box.

Using $\mathrm{Cu} K \alpha$ radiation, equi-inclination multiple film Weissenberg photographs were taken of the zero levels for the crystal rotating about the $a$ and $c$ crystallographic axes and of the zero and first levels for the $b$-axis rotation. The intensities were measured visually using standard calibrated scales and were converted to observed structure factor values in the usual way. No correction was made for absorption. By comparing equivalent planes observed from different rotation axes, the 292 independent reflexions were placed on a common arbitrary scale.

The crystal data are as follows:
$\mathrm{Cu}\left(\mathrm{NO}_{3}\right)_{2} .2 \mathrm{CH}_{3} \mathrm{CN}$, F.W. 269.6. Monoclinic, $a=$ 8.81 (2), $b=7.36$ (2), $c=14.73$ (3) $\AA, \beta=94(1)^{\circ}, U=$ $953 \AA^{3}, \quad D_{m}=1.825 \mathrm{g.cm}^{-3}, \quad Z=4, \quad D_{x}=1.879 \mathrm{~g} . \mathrm{cm}^{-3}$, $\mathrm{Cu} K \alpha$ radiation $(\lambda=1 \cdot 5418 \AA)$. Space group, $P 2_{1} / c$ (no.14).


[^0]:    * Given by numbers only. The first, followed by (2) and a semicolon, are the two atoms lying over and under the metal. E.s.d.'s of Ti-O distances are $\pm 0 \cdot 10 \AA$; for oxygen-oxygen, $\pm 0 \cdot 14 \AA$.

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